

the experimentalist using an LV, the Basset integral term may be neglected since  $C$  is small compared to the other coefficients in Eq. (1). Lastly, the lift force is incorporated into the coefficient  $D$ , only for generality.

### References

- <sup>1</sup>Catalano, G. D., "A Prediction of Particle Behavior via the Basset-Boussinesq-Oseen Equation," *AIAA Journal*, Vol. 23, Oct. 1985, pp. 1627-1628.
- <sup>2</sup>Somerscales, E.F.C., "Measurement of Velocity," *Methods of Experimental Physics, Fluid Dynamics*, Vol. 18A, edited by R. J. Emrich, Academic Press, Orlando, FL, p. 9.

## Comment on "Buckling of Composite Plates with a Free Edge in Edgewise Bending and Compression"

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TABLE 2 of Ref. 1 presented buckling loads for a laminated rectangular plate with two loaded edges simply supported, one side free and the opposite side either simply supported or clamped, for linearly varying edge loads. Since the analysis of Ref. 1 is based on classical plate theory, it is of interest to compare those results with buckling loads that include the effect of transverse shear deformations.

Such buckling loads have been presented for plates with uniformly loaded edges.<sup>2</sup> These were calculated using the shell-of-revolution program FASOR<sup>3</sup> by considering the plate as part of a cylinder of very large radius. The circumferential direction of the cylindrical model corresponds to the longitudinal (loaded) direction of the plate, and the axial length of the cylinder equals the plate width  $b$ . The plate buckling load is obtained by minimizing the model buckling load with respect to the circumferential wave number  $N$  subject to the condition  $N = n\pi R/a$ , where  $n$  is the number of longitudinal half-waves of the plate buckle,  $R$  is the radius of the cylindrical model, and  $a$  is the length of the plate. This procedure is also applicable to plates with linearly varying edge loading.

The dimensions of the laminated plate analyzed in Ref. 1 are  $a = 254$  mm (10 in.) and  $b = 50.8$  mm (2 in.). It is made of 0.127-mm (0.005-in.)-thick tape with the following in-surface elastic properties:<sup>4</sup>  $E_1/E_2 = 10.05$ ,  $G_{12}/E_2 = 0.349$ ,  $\mu_1 = 0.34$ , and  $E_2 = 13.03$  GPa ( $1.89 \times 10^6$  psi), where 1 and 2 signify directions parallel and transverse, respectively, to the fibers and  $\mu_1$  is the major Poisson's ratio. The laminate definition is  $[\pm 45_3/0_3]_s$  (erroneously reported as  $[\pm 45_3/0_3/90_3]_s$  in Ref. 1), thus giving a laminate thickness  $h = 2.286$  mm (0.090 in.). Neglecting the small anisotropic effect, this laminate has the bending stiffness matrix given by Eq. (25) of Ref. 1.

Since the values of the transverse shear moduli  $G_{13}$  and  $G_{23}$  are unavailable, it is assumed that  $G_{13} = G_{23} = G_{12}$ . (For typical unidirectional laminae,  $G_{23} < G_{12}$ .) Again neglecting the small anisotropic effect, FASOR gives the transverse shear stiffness matrix<sup>5</sup>  $[K] = 8.951 \times 10^6$  N/m ( $5.112 \times 10^4$  lb/in.)  $[I]$ , where  $[I]$  is the  $2 \times 2$  unit matrix.

Table 1 compares the classical plate theory results of Ref. 1 with the transverse shear deformation theory results

Table 1 Buckling stress resultants  $N_x b^2/E_2 h^3$  at free edge ( $v = b$ )

$N_x(0)/N_x(b)$	Condition at supported edge ( $v = 0$ )					
	Simply supported			Clamped		
	CPT <sup>a</sup>	FASOR <sup>b</sup>	Diff., %	CPT <sup>a</sup>	FASOR <sup>b</sup>	Diff., %
100	0.101	0.097	4.0	0.219	0.202	8.6
2	2.106	2.018	4.4	4.221	3.954	6.8
1	2.637	2.525	4.4	5.133	4.792	7.1
0.5	3.016	2.886	4.5	5.745	5.357	7.2
0.01	3.507	3.356	4.5	6.491	6.042	7.4
-0.5	4.221	4.036	4.6	7.487	6.952	7.7
-1	5.263	5.027	4.7	8.762	8.121	7.9
-2	10.033	9.579	4.7	12.777	11.744	8.8

<sup>a</sup>Classical plate theory (Ref. 1). <sup>b</sup>Transverse shear deformation theory.

calculated by FASOR. It is noted that the buckling load knockdowns due to transverse shear deformation for this plate are smaller than those reported for uniformly compressed plates in Ref. 2. Aside from differences in laminar properties, lay-up, and edge conditions, this is not surprising since the thickness-to-width ratio for this plate is  $h/b = 0.045$ , which is less than one-half that of the thinnest plate studied in Ref. 2.

### References

- <sup>1</sup>Wang, J. T., Biggers, S. B., and Dickson, J. N., "Buckling of Composite Plates with a Free Edge in Edgewise Bending and Compression," *AIAA Journal*, Vol. 22, March 1984, pp. 394-398.
- <sup>2</sup>Cohen, G. A., "Effect of Transverse Shear Deformation on Anisotropic Plate Buckling," *Journal of Composite Materials*, Vol. 16, July 1982, pp. 301-312.
- <sup>3</sup>Cohen, G. A., "FASOR—A Program for Stress, Buckling and Vibration of Shells of Revolution," *Advances in Engineering Software*, Vol. 3, No. 4, 1981, pp. 155-162.
- <sup>4</sup>Biggers, S. B., Private communication, March 1984.
- <sup>5</sup>Cohen, G. A., "Transverse Shear Stiffness of Laminated Anisotropic Shells," *Computer Methods in Applied Mechanics and Engineering*, Vol. 13, Feb. 1978, pp. 205-220.

## Comment on "Stiffness Matrix Adjustment Using Mode Data"

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THE above paper<sup>1</sup> is a welcome addition to the literatures of structural system identification. The author presents a stiffness matrix adjustment (KMA) procedure that shows some promising results. Nevertheless, we believe that the following comments are appropriate for further developments in the field, especially in relation to real large structures.

Let  $p$  be the number of constraints (equations) and  $q$  the number of unknowns to be identified. Depending on relative values of  $p$  and  $q$ , there are essentially two kinds of system identification.<sup>2,3</sup> For an overdetermined system,  $p > q$ , a least-square solution is sought that minimizes errors of the solution. The procedures given in Refs. 4 and 5 fall into this category. On the other hand, if  $p < q$  (underdetermined system), there are infinitely many sets of solutions satisfying given con-